

Physics 406 Exam 2

Thursday, 4-6-06, 1:20-2:10 PM

1. (25 pts) Suppose a point-like test mass (of mass  $q$ ) is released from rest a distance  $d$  away from approximately two dimensional membrane of mass density  $\sigma$  and infinite extent (i.e. the amount of mass contained the membrane in area  $a$  is  $\sigma a$ ). Assume that the gravitational field generated by the membrane is nonrelativistic and that the membrane is oriented parallel to the  $xy$  plane and is located at  $z = 0$ .
- a) (5 pt) Write down the acceleration of the test mass as a function of  $z$  (the distance to the membrane) using Newton's Law. (Hint: Poisson's equation for gravitation is  $\nabla^2\phi = 4\pi\rho$  where  $\rho$  is the mass density)

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**answer**

Poisson's law gives

$$\nabla^2\phi = 4\pi\sigma\delta(z).$$

Taking a pill box of cross sectional area  $a$ , we have in analogy with electrostatics

$$|\vec{\nabla}\phi|2a = 4\pi\sigma a$$

Hence, we find

$$\frac{d^2z}{dt^2} = 2\pi\sigma$$

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- b) (10 pt) Write down the geodesic equation for the test mass position  $x^\mu(\tau)$  where  $\tau$  is the proper time (leave the answer in terms of  $\Gamma_{\alpha\beta}^\mu$ ).

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**answer**

The geodesic equation is

$$\frac{d^2x^\mu}{d\tau^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0$$

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- c) (10 pt) Assuming that  $\frac{dx^0}{d\tau} = 1$  and non-relativistic speed approximation, compute the Christoffel symbol component  $\Gamma_{tt}^z$  (i.e.  $\Gamma_{00}^3$  in the  $(t, x, y, z)$  coordinates).

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**answer**

The geodesic equation in the nonrelativistic limit is

$$\frac{d^2z}{dt^2} + \Gamma_{00}^3 = 0.$$

Hence, we find

$$\Gamma_{00}^3 = -2\pi\sigma$$

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2. (25 pts) Suppose you are given that at point  $P$  in the coordinate system  $x^\mu = (t, x, y, z)$  that

$$\frac{\partial V^\lambda}{\partial x^\mu} \Big|_P = 0$$

and

$$V^\lambda \Big|_P = (1, 0, 0, 0)$$

where  $V^\lambda$  is a covariant vector under general coordinate transformations. Compute

$$\frac{\partial \bar{V}^0}{\partial \bar{x}^0} \Big|_P$$

(at point  $P$ ) in the coordinate system  $\bar{x}^\mu = (u, v, y, z)$  where

$$u = t - x$$

$$v = t + x$$

assuming that the point  $P$  is located at  $(t, x, y, z) = (1, 0, 0, 0)$ . Hint: Either remember the inhomogeneous transformation rule of  $\partial_\mu V^\lambda$  or even more simply start from the definition of how  $V^\lambda$  transforms under coordinate transformation.

**answer**

Simply expand:

$$\begin{aligned} \frac{\partial \bar{V}^0}{\partial \bar{x}^0} \Big|_P &= \frac{\partial}{\partial \bar{x}^0} \left[ \frac{\partial \bar{x}^0}{\partial x^\lambda} V^\lambda \right] \Big|_P \\ &= \frac{\partial x^\gamma}{\partial \bar{x}^0} \frac{\partial^2 \bar{x}^0}{\partial x^\gamma \partial x^\lambda} V^\lambda \Big|_P + \frac{\partial \bar{x}^0}{\partial x^\lambda} \frac{\partial x^\gamma}{\partial \bar{x}^0} \frac{\partial V^\lambda}{\partial x^\gamma} \Big|_P \\ &= \frac{\partial x^\gamma}{\partial \bar{x}^0} \frac{\partial}{\partial x^\gamma} (1) = 0 \end{aligned}$$

3. (30 pts) Suppose on a two dimensional surface parameterized by  $(x, y)$  (valid for  $x \neq 0$ ) with a metric

$$ds^2 = Q^2 dx^2 + x^4 dy^2,$$

there is a curve  $C$  parameterized by  $(x = x_0, y)$  where  $y$  runs from 0 to 1. (Here,  $Q$  is a constant.) If there is a vector  $\vec{V}$  on the curve at position  $(x = x_0, y = 0)$  with magnitude  $\vec{V}(x = x_0, y = 0) = \hat{y}$  ( $\hat{y}$  is the usual basis unit vector), what is the value of the vector when it is parallel transported along the curve  $C$  to  $(x = x_0, y = 1)$ ? (Hint:  $\Gamma_{2\beta}^2 = \delta_{\beta 1} \frac{2}{x}$  where the notation is as usual, i.e.  $x^1 = x$  and  $x^2 = y$ . Hence, only  $\Gamma_{2\beta}^1$  needs to be computed.)

**answer**

We start with the geodesic equation

$$\frac{dx^\gamma}{d\lambda} \nabla_\gamma V^\alpha = 0.$$

Since we have parameterized the curve in terms of  $y$ , we have

$$\partial_2 V^\alpha + \Gamma_{2\lambda}^\alpha V^\lambda = 0.$$

Writing this out, we have

$$\partial_2 V^1 + \Gamma_{2\lambda}^1 V^\lambda = 0$$

$$\partial_2 V^2 + \Gamma_{2\lambda}^2 V^\lambda = 0$$

Computing

$$\begin{aligned}\Gamma_{2\beta}^1 &= \frac{1}{2} g^{1\lambda} (g_{2\lambda,\beta} + g_{\beta\lambda,2} - g_{2\beta,\lambda}) \\ &= \frac{1}{2} g^{11} (g_{21,\beta} + g_{\beta 1,2} - g_{2\beta,1}) \\ &= \frac{-2x^3}{Q^2} \delta_{\beta 2}.\end{aligned}$$

Hence, our equations turn into

$$\begin{aligned}\partial_2 V^1 - \frac{2x_0^3}{Q^2} V^2 &= 0 \\ \partial_2 V^2 + \frac{2}{x_0} V^1 &= 0\end{aligned}$$

Taking the derivative of the first with respect to  $\partial_2$ , we find

$$\partial_2^2 V^1 - \frac{2x_0^3}{Q^2} \partial_2 V^2 = \partial_2^2 V^1 - \frac{2x_0^3}{Q^2} \left(-\frac{2}{x_0} V^1\right) = \partial_2^2 V^1 + \frac{4x_0^2}{Q^2} V^1 = 0.$$

Since  $V^\lambda$  only depends on  $y$  along the curve of interest, we find

$$V^1 = A \cos\left(\frac{2x_0}{Q} y\right) + B \sin\left(\frac{2x_0}{Q} y\right).$$

Since  $V^1 = 0$  at  $y = 0$ , we have

$$V^1 = B \sin\left(\frac{2x_0}{Q} y\right).$$

Since  $V^2 = 1$  at  $y = 0$ , we have the second boundary condition giving

$$V^2 = \frac{Q^2}{2x_0^3} \partial_2 V^1 = \frac{Q}{x_0^2} B \cos\left(\frac{2x_0}{Q} y\right)|_{y=0} = 1.$$

Hence,

$$\begin{aligned}V^1 &= \frac{x_0^2}{Q} \sin\left(\frac{2x_0}{Q} y\right) \\ V^2 &= \cos\left(\frac{2x_0}{Q} y\right)\end{aligned}$$

4. (20 pts) Using the doubly contracted Bianchi identity

$$\nabla_\alpha (R^{\alpha\beta} - \frac{1}{2} g^{\alpha\beta} R) = 0,$$

find the partial differential equation satisfied by the function  $K(x)$

$$R_{\alpha\gamma} = K(x) g_{\alpha\gamma}.$$

**answer**

$$R_{\alpha\gamma} = K(x)g_{\alpha\gamma}$$

$$R = 4K(x)$$

Since the covariant derivative of the metric vanishes, we have

$$\nabla_{\alpha}[R^{\alpha\gamma} - \frac{1}{2}g^{\alpha\gamma}R] = 0$$

implying

$$K_{,\alpha}g^{\alpha\gamma} - 2g^{\alpha\gamma}K_{,\alpha} = -g^{\alpha\gamma}K_{,\alpha} = 0$$

Hence,  $K$  must be a constant.

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