

# Longitudinal Waves

Speed of propagation  
Sound Waves



Air, Liquid (Tsunami)  
Solid (Earthquakes)

Longitudinal + Transverse

Hear  $20\text{Hz} \Rightarrow 20\text{kHz}$   
 Below Audible infrasonic  
 Above Audible Ultrasonic

For each type of wave  
 Speed  $C = \frac{\text{Elastic Modulus}}{\text{density}}$   
 see serway 12.4

COMPRESSION WAVES  
 $B = \text{Bulk Modulus}$  } = density  
 $C = \sqrt{\frac{B}{\rho}}$  Primary wave of earthquake

Secondary earthquake waves involve shear = transverse waves  $\Rightarrow$  use shear modulus

$$C = \sqrt{\frac{S}{\rho}} = S$$

Sound Waves  
 Bulk Modulus for Gas has different values depending on compression type

Compression at constant temperature (isothermal) or without heat transfer (Adiabatic). For Sound Adiabatic

$$C = \sqrt{B(\text{adiabatic})}$$

$$= \sqrt{\gamma B(\text{isothermal})}$$

$$\gamma = \frac{C_p}{C_v} \quad (\text{serway } 21.2)$$

$C_p$  = specific heat at const. press.

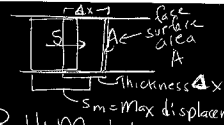
$C_v$  = specific heat at constant volume

Air:  $\gamma = 1.4$   $P =$

$$C = \sqrt{\frac{\gamma P}{\rho}} \quad \text{Pressure}$$

$$V = 344 \text{ m/s} = 1139 \text{ ft/s}$$

Sound waves in air: elements of air



Bulk Modulus

$$B = -\frac{\Delta P}{\Delta V/V}$$

Distance btw front and back changes from  $\Delta x$  by  $\Delta s$  due to compression.

Volume was  $V = A\Delta x$

$$\Delta V = A \Delta S \Rightarrow$$

$$\Delta P = -B \frac{\Delta V}{V} =$$

$$-B \frac{A \Delta S}{A \Delta x} = -B \frac{\Delta S}{\Delta x}$$

$$\lim_{\Delta x \rightarrow 0} \Rightarrow dP = -B \frac{\partial S}{\partial x}$$

Oscillating Displacement  
 $A = \text{Amplitude} = S_m$

$$S = S_m \cos(kx - \omega t)$$

$$dP = -B \frac{\partial}{\partial x} (S_m \cos(kx - \omega t))$$

$$= +Bk S_m \sin(kx - \omega t)$$

$$\Delta P_m = Bk S_m, C = \sqrt{\frac{B}{\rho}}$$

$$C^2 = \frac{B}{\rho} \Rightarrow B = C^2 \rho$$

$$\Delta P_m = (C^2 \rho k) S_m$$

$$C = \frac{\omega}{k} \Rightarrow Ck = \omega$$

$$\Delta P_m = (C^2 \rho \omega) S_m$$

Thin slice of air thickness  $dx$ , mass  $dm$ , area  $A$  oscillates as sound passes

Kinetic Energy:  $dK = \frac{1}{2} dm v_s^2$

$v_s =$  speed of oscillating air element, not wave speed

$$v_s = \frac{\partial S}{\partial t} = -\omega S_m \sin(kx - \omega t)$$

$$dm = \rho A dx \Rightarrow dK = \frac{1}{2} dm v_s^2$$

$$dk = \frac{1}{2} (\rho A dx) (-\omega S_m)^2 \sin^2(kx - \omega t)$$

Power is  $dk/dt$ ,  $c \equiv \frac{dx}{dt}$

$$\frac{dk}{dt} = \frac{1}{2} \rho A c \omega^2 S_m^2 \sin^2(kx - \omega t)$$

Average over time:  $\left(\frac{dk}{dt}\right) = \frac{1}{2} \rho A c \omega^2 S_m^2 \sin^2(kx - \omega t)$   
 $= \frac{1}{4} \rho A c \omega^2 S_m^2$

Like a spring on average, energy is  $\frac{1}{2}$  potential +  $\frac{1}{2}$  kinetic energy

$$\frac{dE}{dt} = 2 \frac{dk}{dt} = \frac{1}{2} \rho A c \omega^2 S_m^2$$

Intensity =  $\frac{\text{Power}}{\text{Area}} = \frac{dE/dt}{A_{\text{area}}} = \frac{1}{2} \rho c \omega^2 S_m^2$   
 ( $S_m = A = \text{amplitude}$ )

$$I = \frac{1}{2} \rho c \omega^2 A^2 \quad \text{where } A = \text{amplitude}$$

$$\Delta P_m = \rho c \omega S_m \Rightarrow I = \frac{1}{2} \rho c \omega^2 S_m^2$$

Can be written:  $I = \frac{\Delta P_m^2}{2 \rho c}$   
 Sounds vary in intensity by factor of  $10^{12}$

$$\text{Decibel} = \text{db} = 10 \log \frac{I}{I_0}, \quad I_0 = 10^{-12} \frac{\text{watts}}{\text{m}^2}$$

Sound from moving sources

# Doppler Effect

Detector moves towards stationary source at  $v_D$



If  $v_D$  were zero in time  $t$   $ct/\lambda$  waves pass

Because detector moves additional  $v_D t/\lambda$  waves pass. So instead of hearing  $f = \frac{c}{\lambda} = \frac{ct}{\lambda t}$

$$\text{hear } f' = \frac{ct}{\lambda} + \frac{v_D t}{\lambda}$$

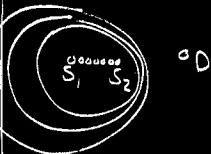
$$= \frac{c + v_D}{\lambda} = \frac{c + v_D}{c/f} = f \left( \frac{c + v_D}{c} \right)$$

$$f' = f \left( \frac{c + v_D}{c} \right)$$

If moving away, waves chase detector, fewer arrive:

$$f' = f \left( \frac{c - v_D}{c} \right)$$

Supposed detector at rest.  $v_D = 0$   
Source moves at  $v_s$



Let  $T = 1/f = \text{Period}$   
 During one period  $T$   
 Source moves  $V_s T = \frac{V_s}{f}$   
 Towards detector.

Wavelength  $\lambda$  &  $f$   
 arriving sound is not  
 $\lambda = \frac{c}{f}$  but it is

$$\lambda' = \frac{c - V_s}{f}$$

We hear  $f'$

$$f' = \frac{c}{\lambda'} = \frac{f}{1 - \frac{V_s}{c}}$$

Move away:  $V_s \rightarrow -V_s$   
 $f' = f \left( \frac{c}{c + V_s} \right)$

If both source & detector  
 move

$$f' = f \left( \frac{c \pm V_p}{c \mp V_s} \right)$$

Towards  
away