

HOMEWORK 4 SOLUTIONS

3-3b) (a) $E = hf = \frac{hc}{\lambda} = \frac{1240 \text{ eV} \cdot \text{nm}}{0.0711 \text{ nm}} = [1.75 \times 10^4 \text{ eV}] = 17.5 \text{ keV}$

(b) For $\theta = 180^\circ \cos\theta = -1$ so we get

$$\begin{aligned}\lambda_2 &= \lambda_1 + \frac{h}{mc}(1-\cos\theta) \\ &= 0.0711 + [(1240 \text{ eV} \cdot \text{nm}) / (5.11 \times 10^5 \text{ eV})] (2) \\ \boxed{\lambda_2 = 0.0760 \text{ nm}}\end{aligned}$$

(c) The corresponding photon energy is

$$E_2 = \frac{hc}{\lambda_2} = \frac{1240 \text{ eV} \cdot \text{nm}}{0.0760 \text{ nm}} = [1.63 \times 10^4 \text{ eV}] = 16.3 \text{ keV}$$

(d) Total energy is conserved, so the electrons K.E. is

$$\begin{aligned}E &= 1.75 \times 10^4 \text{ eV} - 1.63 \times 10^4 \text{ eV} \\ \boxed{E = 1.130 \text{ eV}} &= 1.130 \text{ keV}\end{aligned}$$

3-4g) To get useful information from the data it's best to make a graph like Fig 3-13. The stopping potential is given by

$$\begin{aligned}q \cdot V_0 &= hf - \phi \\ &= hc/\lambda - \phi\end{aligned}$$

so V_0 is a linear function of f or of $1/\lambda$. For each wavelength lets find the frequency using $\lambda f = c$

$$\lambda = 253.5 \text{ nm} \quad f = 1.183 \times 10^{15} \text{ /s}$$

312.5	0.960	"
365.0	0.822	"
404.7	0.741	"
433.9	0.691	"

(a) The graph is on the next page. To get the work function draw a line through the points and extrapolate to $f=0$. At $f=0$ we have $qV_0 = -\phi \Rightarrow I$ get

$$\boxed{\phi \approx 2.3 \text{ eV}}$$

(b) We get Planck's constant from the slope. I get

$$h = \frac{1.2 \text{ eV}}{0.3 \times 10^{15} \text{ s}} \Rightarrow h = 4.0 \times 10^{-15} \text{ eV.s}$$

(c) The work function for lead is 4.14 eV, which is $(4.14 - 2.3) eV
 $= 1.84 \text{ eV}$ greater than lithium \Rightarrow the electron energies would be smaller by 1.84 eV. We get emission only for 253.5 nm.$

8-27) The average energy is given by Eq. 8-84

$$\langle E \rangle = \frac{hf}{e^{hf/kT} - 1} = kT \left[\frac{hf/kT}{e^{hf/kT} - 1} \right]$$

(a) $T = 10 \text{ hf}/k$
 $\Rightarrow hf/kT = 10$

$$\langle E \rangle = (kT) \frac{0.1}{e^{0.1} - 1}$$

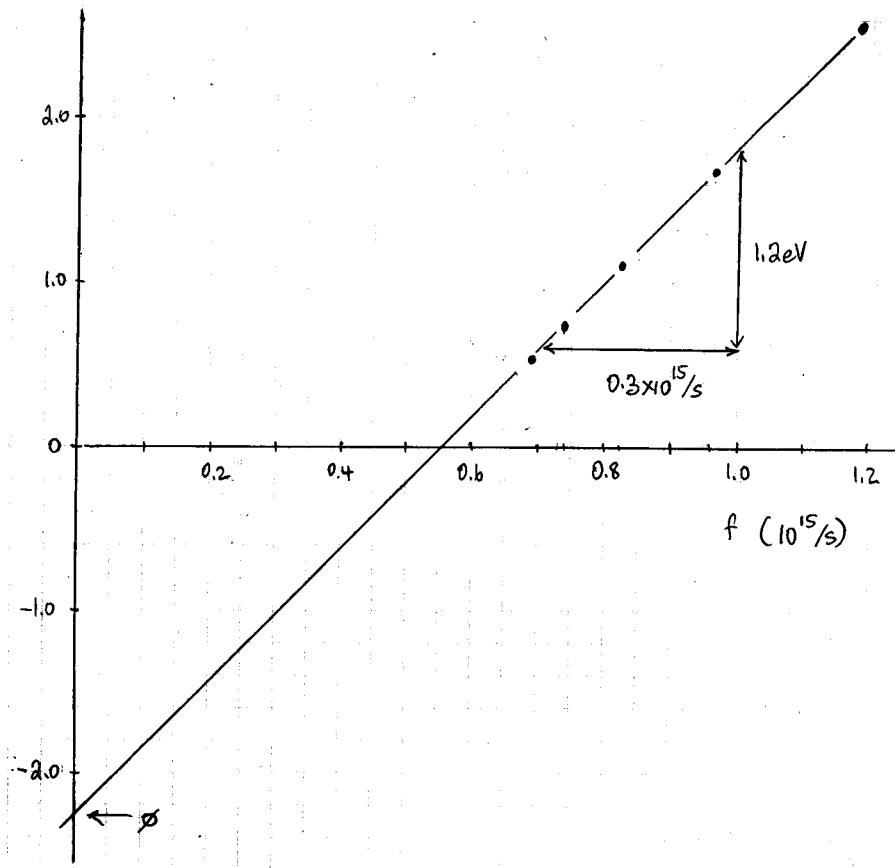
$\langle E \rangle = 0.951 kT$ which is close to the equipartition value.

(b) $T = hf/k \Rightarrow hf/kT = 1$

(c) $T = 0.1 \frac{hf}{k} \Rightarrow \frac{hf}{kT} = 10$

$$\langle E \rangle = kT \frac{1}{e^1 - 1} = \langle E \rangle = 0.582 kT$$

$$\langle E \rangle = kT \frac{10}{e^{10} - 1} \Rightarrow \langle E \rangle = 0.000454 kT$$



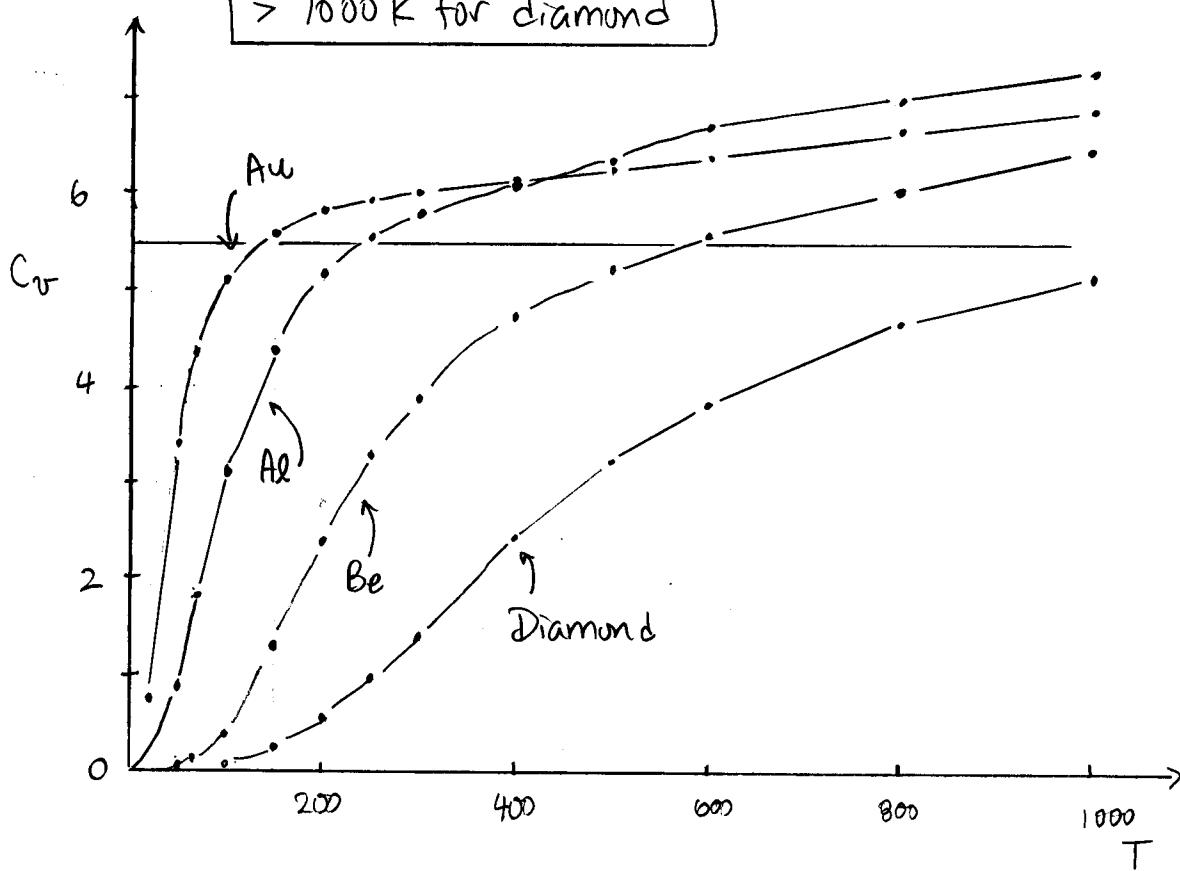
8-31) From Eq. 8-86 we can find C_V at temperature $T = T_E$.

$$\begin{aligned} C_V &= 3R \left(\frac{T_E}{T}\right)^2 \frac{e^{T_E/T}}{(e^{T_E/T}-1)^2} = 3R(1)^2 \frac{e^1}{(e^1-1)^2} \\ &= 3R(0.92) \\ &= 3(1.987 \text{ cal/mole}\cdot\text{K})(0.92) = 5.49 \text{ Cal/mole}\cdot\text{K} \end{aligned}$$

So T_E is roughly the temperature where $C_V = 5.5 \text{ cal}$.

From the graph I get

$T_E = 140 \text{ K}$ for gold
 250 K for Al
 580 K for Be
 $> 1000 \text{ K}$ for diamond



4-5) The wavelengths are supposed to be given by the formula

$$\frac{1}{\lambda} = R \left(\frac{1}{m^2} - \frac{1}{n^2}\right) \quad n > m$$

where $R = 1.096776 \times 10^7 / \text{m} = 0.01096776 / \text{nm}$.

Trying random values I get

$$\lambda = 7460$$

$$m=5, n=6$$

$$\lambda = 4654$$

$$m=5, n=7$$

$$\lambda = 3741$$

$$m=5, n=8$$

There is no way to get 4103 nm

4-14) From 4-19 $a_0 = \frac{\hbar^2}{mke^2}$. In the books notation ($k = \frac{1}{4\pi\epsilon_0}$) we have $\alpha = ke^2/mc$ so we can write

$$a_0 = \frac{\hbar^2 c^2}{mc^2 ke^2} = \left(\frac{\hbar c}{mc^2}\right) \left(\frac{\hbar c}{ke^2}\right) = \left(\frac{\hbar c}{mc^2}\right) \frac{1}{\alpha} = \frac{\hbar}{2\pi mc} \frac{1}{\alpha}$$

So

$$a_0 = \frac{\lambda_c}{2\pi} \frac{1}{\alpha} \quad \lambda_c = \frac{\hbar}{mc}$$

From 4-20

$$E_1 = -\frac{mk^2 z^2 e^4}{2\hbar^2 n^2}$$

So the magnitude of E_1 for hydrogen ($z=1$) is

$$|E_1| = \frac{1}{2} mc^2 \frac{k^2 e^4}{\hbar^2 c^2} \left(\frac{1}{1^2}\right) \Rightarrow E_1 = \frac{1}{2} \alpha^2 mc^2$$

4-15) The wavelengths are given by $\frac{1}{\lambda} = R \left[\frac{1}{n^2} - \frac{1}{m^2} \right]$ where $m > n$ (so m is the initial state and n is final state). For the Lyman Series $n=1$. I get the following results

$m=2$	$\lambda = 121.57$
$m=3$	$\lambda = 102.57$
$m=4$	$\lambda = 97.25$

The series limit (setting $m \rightarrow \infty$) is $\lambda = 91.18$ nm

These are all well outside the visible region.

4-41) From 4-14 we had $a_0 = \frac{\hbar}{mc} \frac{1}{\alpha}$ so in general $r_n = \frac{\hbar}{mc} \frac{1}{\alpha} n^2$. But $|E| = mr^2 = n\hbar \Rightarrow$

$$r = n\hbar/mr = (n\hbar/m) \left(\frac{mc}{\hbar}\right) \frac{\alpha}{n^2} \Rightarrow r = \alpha c/n$$