

# *Black Hole Attractors and Superconformal Quantum Mechanics*

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- *Brief review:*
  - *BPS black holes in  $\mathcal{N} = 2$  supergravity*
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  - *Short multiplets*

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- *BH entropy from D2-brane wrapping horizon.*

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  - $h_{2,1} + 1$  hypermultiplets

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- $F_A = \partial_A F = \partial_A \left[ \frac{D_{abc} X^a X^b X^c}{X^0} + \dots \right]$

# *BPS Black Holes*

- *Graviphoton:*  $X^A \mathcal{G}_A - F_A \mathcal{F}^A$
- *Mass:*  $Z_\infty = (p^A F_A - q_A X^A) |_\infty$

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  - *Depends on moduli at infinity.*
- *Dependence is forgotten near the horizon*
- *$X^A(r) \rightarrow X_{fixed}^A(p^A, q_A)$  as  $r \rightarrow 0$* 
  - *Scalars flow to an attractor point*

# BPS Black Holes

- $X_{fixed}^A$  extremizes  $|Z| = |p^A F_A - q_A X^A|$
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- $\Re C X_{fixed}^A = p^A$        $\Re C F_{A fixed} = q_A$
- $S_{BH}(p, q) = \pi |Z|_{fixed}^2$
- $\exp S_{BH}(p, q) \simeq$  Number of microstates of black hole

# Entropy from microstates

- Charges  $\begin{pmatrix} p^A \\ q_A \end{pmatrix}$  from *D-branes in Calabi-Yau*  
 $(p^0, p^a, q_a, q_0)$        $(D6, D4, D2, D0)$

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- Counting of bound states of  $D4, D2, D0$ 
  - $S_{BH}(p, q) \simeq 2\pi \sqrt{D\tilde{q}_0}$
  - $D = D_{abc}p^a p^b p^c$        $\tilde{q}_0 = q_0 + \frac{1}{12}D^{ab}q_a q_b$

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  - Different computations:  $M$ -theory, others.
  - $\prod \frac{1}{(1-x^n)^{6D}}$

# *Probing near the horizon*

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$$u^A q_A - v_A p^A > 0$$
- *Many ways to probe near horizon geometry.*

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## *Superconformal QM of probe D-branes*

- *Black Hole Entropy:*
  - *microstates as bound states in near-horizon geometry*
- *AdS<sub>2</sub>/CFT<sub>1</sub> correspondence:*
  - *Superconformal Matrix Quantum Mechanics?*

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- *Charges*  $\begin{pmatrix} u^A \\ v_A \end{pmatrix}$   $m = |u^A F_A - v_A X^A|$

- $q_m = p^A v_A - q_A u^A \quad q_e^2 + q_m^2 = m^2$

# Hamiltonian

- *Hamiltonian to quadratic order ( $\sigma = \xi^2$ )*

- $$\frac{1}{8mR} \hat{P}_\xi^2 + \frac{1}{2mR\hat{\xi}^2} \hat{L}_{S_2}^2 + \frac{1}{2mR\hat{\xi}^2} \hat{\Delta}_{CY}$$

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- $\mathcal{N} = 4$  *superconformal symmetry*  $SU(1, 1|2)$
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- $SU(2)$  *R-symmetry*  $\hat{L}_i^{tot} = \hat{L}_i^{S_2} + \hat{L}_i^{spin}$
- $\hat{L}_{\alpha\beta}^{tot} = \hat{L}_i^{tot} \sigma_{\alpha\beta}^i$

# Fermions

- Fermionic partners:  $D0$  brane has 16 in  $10d$  spinor  $\Theta$
- $\Theta = \lambda_\alpha u + \eta_\alpha^a \Gamma_a u + \bar{\eta}_\alpha^{\bar{a}} \Gamma_{\bar{a}} \bar{u} + \bar{\lambda}_\alpha \bar{u}$

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- $\{q_\alpha^A, q_\beta^B\} = \hat{\Delta}_{CY} \epsilon_{\alpha\beta} \epsilon^{AB}$
- $R$ -symmetry  $SU(2)$  acts as Lefschetz action on  $(p, q)$ -forms  $(J \wedge, (p - q)/2, \iota(J))$

# Fermionic symmetries

- $\{\hat{S}_\alpha^A, \hat{S}_\beta^B\} = 2\epsilon^{AB}\epsilon_{\alpha\beta}\hat{K}$
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- $\{\hat{S}_\alpha^A, \hat{Q}_\beta^B\} = \epsilon^{AB}(\epsilon_{\alpha\beta}\hat{D} + 2i\hat{L}_{\alpha\beta}^{tot})$

# Supersymmetry generators

- $\hat{Q}_\alpha^A = \frac{1}{\hat{\xi}} \hat{q}_{\alpha, CY}^A +$   
 $(\frac{1}{2} \hat{P}_\xi \lambda_\alpha^A - \frac{i}{\hat{\xi}} (\hat{L}_{\alpha\beta}^{S_2} + \hat{L}_{\alpha\beta}^{CY} + \frac{1}{2} L_{\alpha\beta}^\lambda) \lambda^{A\beta}$

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- $\{\hat{Q}_\alpha^A, \hat{Q}_\beta^B\} = 2\epsilon^{AB} \epsilon_{\alpha\beta} \hat{H}$
- *Nontrivial cancellations at work.*
- $\hat{H} = \frac{1}{8mR} \hat{P}_\xi^2 + \frac{1}{2mR\hat{\xi}^2} \hat{L}_{tot}^2 + \frac{m}{R\hat{\xi}^2} \hat{\Delta}_{CY} + \dots$

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- *In short multiplet also  $G_{+, -\frac{1}{2}}^A$*

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- *State counting depends only on CY topology*

# *Impasse*

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- $F_4 = \omega_{S_2} \wedge \sum p^A \alpha_A$
- *Difficult to write Matrix SCQM*

# *D2-branes wrapping $S^2$*

- *No flat directions from gauge field*
- *Drop  $L^{S^2}$  part from SCQM*
- *Magnetic field on CY :  $F_{a\bar{b}} = \sum p^A \alpha_A = J_{a\bar{b}}$*

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- *Magnetic field changes anticommutator of  $q_\alpha^A$*
- *$\{q_\alpha^A, q_\beta^B\} = \hat{H}^{CY} \epsilon_{\alpha\beta} \epsilon^{AB} + c' L_{\alpha\beta}^{CY} \sigma_1^{AB}$*
- *Spurious term in  $\{Q_\alpha^A, Q_\beta^B\}$*

# Corrected supercharge

- $$\hat{Q}_\alpha^A = \left( \frac{1}{2} \hat{P}_\xi \lambda_\alpha^A - \frac{i}{\hat{\xi}} (\hat{L}_{\alpha\beta}^{CY} + \frac{1}{2} L_{\alpha\beta}^\lambda) \right) \lambda^{A\beta} +$$
$$+ c(\sigma_3)^A_B \lambda_\alpha^B + \frac{1}{\hat{\xi}} \hat{q}_{\alpha, CY}^A$$

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- *Extra potential term in bosonic Hamiltonian from mass of  $D2$*
- $\{\hat{S}_\alpha^A, \hat{Q}_\beta^B\} = \dots + c(\sigma_1)^{AB}\epsilon_{\alpha\beta}$
- *Central extension from charge of  $D2$  on sphere*

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- *Ground states from Landau levels on Calabi Yau.*
- *Leading order  $4D$  bosonic ground states,  $4D$  fermionic*
- $D = D_{abc}p^a p^b p^c$

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- *Same counting as usual: different ways to partition  $q_0$  among  $D2$  branes of  $6D$  kinds*
- *Asymptotic number of states  $\exp(2\pi\sqrt{Dq_0})$*

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- *Why is it interesting?*
  - $Z_{BH} = \sum_q \Omega(p, q) e^{-q_A \phi^A} \simeq |Z|_{top}^2$
  - $\Omega(p, q)$  is number of bound states: open string theory, ignore attractor eqn.
  - $|Z|_{top}^2$  is closed string topological partition function at the attractor point!
- *Natural to seek explanation in attractor geometry.*

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