

Lecture 5 Addendum

① $W_{\mu} = \frac{1}{2} \epsilon_{\mu\nu\rho\sigma} P^{\nu} M^{\rho\sigma}$

↓

$$M^{\rho\sigma} = \underbrace{S^{\rho\sigma}}_{\text{spin}} + i(x^{\rho} \partial^{\sigma} - x^{\sigma} \partial^{\rho})_{\text{orbital}}$$

Orbital part drops out

$$W_{\mu} = \frac{1}{2} \epsilon_{\mu\nu\rho\sigma} P^{\nu} S^{\rho\sigma}$$

For massive states $W^2 = -m^2 s(s+1)$ intrinsic spin

② For massless states $W^2 = P^2 = 0$ & $W_{\mu} P^{\mu} = 0$

$\Rightarrow W_{\mu} = \lambda P_{\mu}$

Consider W_0 component

$$\lambda = \frac{\vec{P} \cdot \vec{S}}{P_0} = \hat{P} \cdot \vec{S}$$

def of helicity

λ has magnitude = $|\vec{S}|$ and is real

$\therefore \lambda = \pm S$

③ Recall in two-component form

$$\{ Q_\alpha, \bar{Q}_{\dot{\beta}} \} = 2 \sigma^u_{\alpha\dot{\beta}} P_u$$

$$\{ Q_\alpha, Q_\beta \} = \{ \bar{Q}_{\dot{\alpha}}, \bar{Q}_{\dot{\beta}} \} = 0$$

Up to normalization, Q_α and $\bar{Q}_{\dot{\beta}}$ are the annihilation & creation operators of Clifford algebra

Last time, we constructed a state

$$|p, \lambda'\rangle = Q_1 Q_2 |p, \lambda\rangle$$

This is the Clifford vacuum.

For the same reason that $\bar{Q}_{\dot{i}}$ changes the helicity by $\frac{1}{2}$, Q_1 & Q_2 also changes the helicity by $\frac{1}{2}$ unit.

This follows from

two component form

$$\left\{ \begin{aligned} [M_{uv}, Q_\alpha] &= -\frac{1}{2} (\sigma_{uv})_{\alpha\beta} Q_\beta \\ [M_{uv}, \bar{Q}_{\dot{\alpha}}] &= -\frac{1}{2} (\bar{\sigma}_{uv})_{\dot{\alpha}\dot{\beta}} \bar{Q}_{\dot{\beta}} \end{aligned} \right.$$

massless state

$$[W_0, \bar{Q}^{\dot{\alpha}}] |p, \lambda\rangle = -\frac{1}{2} P_0 (\sigma^3 Q)^{\dot{\alpha}} |p, \lambda\rangle$$

etc

④ Particles with spin $> 1/2$ related to gauge sym
to decouple ghosts in S-matrix

Spin 1 : Yang-Mills

Spin 2 : general coordinate inv.

Spin $3/2$: local SUSY

Given a Spin $3/2$ field (Rarita-Schwinger field)
construct a current

$$J_\mu = \epsilon_{\mu\nu\rho\sigma} \gamma^\nu \partial^\rho \psi^\sigma$$

↑
Spin $3/2$ field

suppress
spinorial
index

$$\partial^\mu J_\mu = 0$$

Introduce a term $\mathcal{L}' = -i \bar{\psi}_\mu J^\mu$
to cancel extra terms since we promote
SUSY to local $\alpha \rightarrow \alpha(x)$

$$\text{But } \delta_{\alpha(x)} \mathcal{L}' = i k \bar{\psi}_\mu \gamma_\nu T^{\mu\nu} \alpha + \dots$$

↑
stress tensor

Need $\mathcal{L}_{\text{grav}} = -g_{\mu\nu} T^{\mu\nu}$ to cancel these
terms

\Rightarrow SUGRA, We need to include spin 2 field
in multiplet,