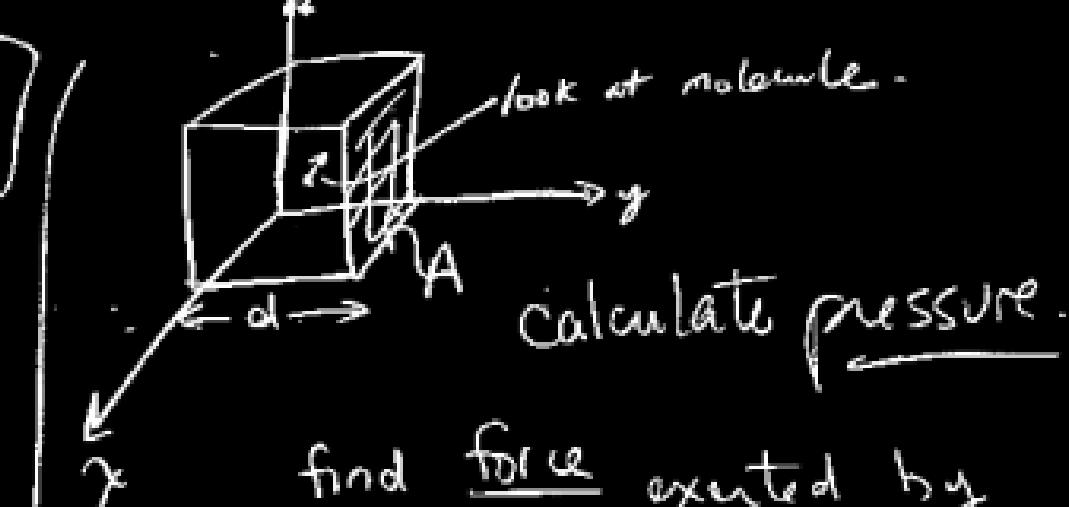
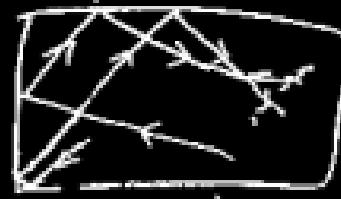


# ch 21 Kinetic theory of gases

## Assumptions:

- 1) large # of molecules  
large separation between molecules
- 2) molecules obey Newton's laws  
as a whole they move randomly  
"molecular chaos"
- 3) molecules interact via <sup>only</sup> short-range  
forces during elastic collisions
- 4) molecules make elastic collisions with walls
- 5) all molecules in gas are identical



calculate pressure



find force exerted by  
molecules on box

momentum change of  $i^{th}$   
particle when it collides  
with wall is

$$\Delta p = -2Mv_{xi} \lambda$$



$$\frac{\text{change in momentum}}{\text{Unit time}} = \bar{F}_x \leftarrow \begin{matrix} \text{for} \\ \text{particles} \\ \text{exerted by} \\ \text{wall} \end{matrix}$$

$$\frac{(-2mV_{xi})}{(2d/V_{xi})} \stackrel{\text{time between}}{\curvearrowright} \text{collisions} = \frac{(2d)}{V_x}$$

$$\text{average force} = -\frac{mV_{xi}^2}{d}$$

on 1 particle  
colliding with  
right wall

$$\text{average force exerted} = \frac{mV_{xi}^2}{d}$$

by  $i^{th}$  particle on wall

If lots of particles  
average force is the force (basically)

$$\text{total force} = \frac{m}{d} \sum_i V_{xi}^2$$

$$= \frac{m}{d} N \overline{V_{xi}^2} \quad \begin{matrix} \text{average of} \\ V_{xi}^2 \end{matrix}$$

To calculate  $\overline{V_{xi}^2}$ , note

$$\overline{V^2} = \overline{V_x^2} + \overline{V_y^2} + \overline{V_z^2}$$

$$+ \overline{V_x^2} = \overline{V_y^2} = \overline{V_z^2}$$

$$\Rightarrow \overline{V_x^2} = \frac{1}{3} \overline{V^2}$$

$$F = \frac{m}{d} \frac{N}{3} \overline{V^2}$$

$$P = F/A = \left(\frac{2}{3}\right) \left(\frac{N}{dA}\right) \left(\frac{1}{2} m \overline{V^2}\right)$$

$$Ad = V_1 \infty$$

$$PV = \frac{2}{3} N \left( \frac{1}{2} m \bar{v^2} \right)$$

compare to  $PV = nRT$

$$\Rightarrow T = \frac{2}{3k_B} \left( \frac{1}{2} m \bar{v^2} \right)$$

relates  $T$  to  $\bar{KE}$

$$\frac{1}{2} m \bar{v^2} = \frac{3}{2} k_B T \quad (k_B = \frac{R^{gas}}{N})$$

$$\text{note } \bar{v^2} = \bar{v_x^2} + \bar{v_y^2} + \bar{v_z^2}$$

$\frac{1}{2} k_B T$  per translational  
degree of freedom.

ideal gas,

$E_{int}$  depends only on  $T$

(not true for non-ideal gas)

related  $\bar{v^2}$  to  $T$

root mean square speed  $v_{rms} = \sqrt{\bar{v^2}}$

$$v_{rms} = \sqrt{\bar{v^2}} = \sqrt{\frac{3k_B T}{m}} = \sqrt{\frac{3RT}{M}}$$

↑  
molecule mass in kg.

note:  
lighter molecules move faster.

calculate specific heat for  
ideal gas.

last time:

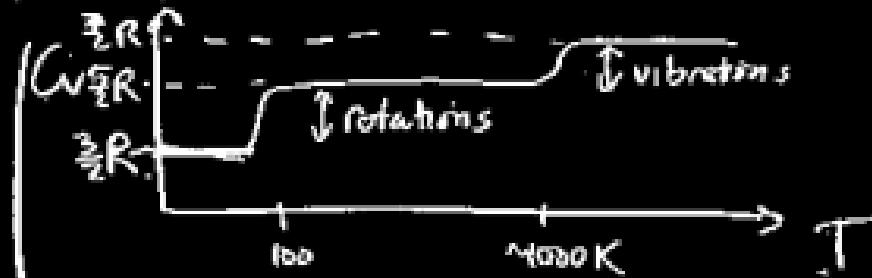
$$C_V = \frac{dE_{int}}{dT}$$

for monoatomic atoms,  $E_{int} = \frac{3}{2}nRT$

$$\Rightarrow C_V = \frac{3}{2}R$$

now consider diatomic ideal gas

$N_2, O_2$  translations + 2 rotational d.o.f freedom  
+ 2 vibrational degrees of freedom



at room temperature, air has  $C_V \sim \frac{5}{2}R$

translations - 3 d.o.f  
rotations - 2 d.o.f

$$\text{internal energy } \frac{3}{2}nRT \rightarrow \underline{\underline{\frac{5}{2}nRT}} \rightarrow \underline{\underline{\frac{7}{2}nRT}}$$

Specific heat at constant pressure



1<sup>st</sup> law of thermo

$$dE_{int} = dQ + dW$$

$$dQ = dE_{int} - dW$$

fixed P:

$$nC_p \Delta T = nC_v \Delta T - (-P \Delta V)$$

For ideal gas,  $PV = nRT$

for this path,  $P \Delta V = nR \Delta T$

$$\Delta V = \frac{nR}{P} \Delta T$$

So  $nC_p \Delta T = nC_v \Delta T + P \left( \frac{nR}{P} \right) \Delta T$

$$\therefore C_p = C_v + R$$

$$\gamma = \frac{C_p}{C_v} = \frac{C_v + R}{C_v} = 1 + \frac{R}{C_v}$$

for monatomic ideal gas

$$C_v = \frac{3}{2} R$$

$$\gamma = \frac{5/2 R}{3/2 R} = \frac{5}{3}$$

for diatomic ideal gas,

$$\gamma = \frac{7/2 R}{5/2 R} = 7/5$$

So far have related  $\frac{T}{T_0}$  to  $\frac{1}{2mV^2}$   
now discuss probability distribution

Boltzmann distribution

Number density =  $\frac{\# \text{ of molecules}}{\text{volume}}$  with energy between  $E$  and  $E + dE$   $\equiv n_v(E) = n_0(v) e^{-E/k_B T}$

