

$$dE_{int} = dQ + dW \quad [-PdV]$$

adiabatic expansion  
of ideal gas

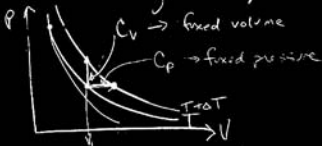
$$dQ = 0 \Rightarrow dE_{int} = dW$$

recall specific heat measures  
amount of heat needed to change  
temp by  $\Delta T$

$$\Delta Q = nC\Delta T$$



specific heat depends on  
how you change  $T$



look at  $C_v$  :  
at constant  $V$   $\rightarrow [v=0]$   
 $dE_{int} = dQ + dW$   
 $dE_{int} = nC_v dT$

$$\text{So } nC_v dT = dE_{\text{int}}$$

and

$$nC_v dT = -PdV$$

now eliminate  $dT$  in favor  
of  $P \leftarrow V$

$$PV = nRT$$

$$PdV + VdP = nR dT$$

$$\Rightarrow dT = \frac{1}{nR} (PdV + VdP)$$

$$\frac{nC_v}{nR} (PdV + VdP) = -PdV$$

$$VdP = -(1 + R/C_v) PdV$$

$$\frac{dP}{P} = -\left(1 + \frac{R}{C_v}\right) \frac{dV}{V}$$

$$\ln P = C_1 - \left(1 + \frac{R}{C_v}\right) \ln V$$

exponentiate:

$$P = C V^{-\left(1 + \frac{R}{C_v}\right)}$$

$$\Rightarrow PV^{\left(1 + \frac{R}{C_v}\right)} = \text{const}$$

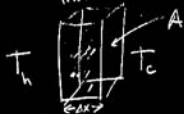
will show later that for an ideal gas,

$$1 + R/C_v = C_p/C_v$$

define  $\gamma = C_p/C_v$

along adiabat  $PV^\gamma = C$

Energy transfer -  
thermal conductivity



find empirically that heat  $Q$   
transferred in a time  $\Delta t$   
from hot side to cold side obeys

$$\frac{Q}{\Delta t} \propto A (T_h - T_c) / \Delta x$$

define  $\mathcal{P} = \frac{Q}{\Delta t}$  = power transferred  
(watts if  $Q$  is in J  
and  $\Delta t$  is in s)

define thermal conductivity  $k$  by

$$\mathcal{P} = k A \left| \frac{dT}{dx} \right|$$

units of  $k$  are  $\frac{W}{m \cdot K}$

for slabs,  $\mathcal{P} = k A \frac{T_h - T_c}{\Delta x}$

high  $k$  = good thermal conductor  
low  $k$  = poor thermal conductor

relate R value for insulation  
to thermal conductivity:

$$\text{define } R = \frac{\Delta x}{k}$$

bigger R  $\Rightarrow$  better insulator

typical units for R are "English"  
 $\text{ft}^2 \cdot ^\circ\text{F} \cdot \text{h} / \text{Btu}$



two slabs in "series"

heat cannot accumulate at

interface, so  $Q_1 = k_1 A \left( \frac{T_h - T_i}{L_1} \right) = Q_2 = k_2 A \left( \frac{T_i - T_c}{L_2} \right) \Rightarrow Q = \frac{A(T_h - T_c)}{\left( L_1/k_1 + L_2/k_2 \right)}$

$$\frac{k_1}{L_1} (T_h - T_i) = \frac{k_2}{L_2} (T_i - T_c)$$

find  $T_i$

$$T_i (k_1 L_2 + k_2 L_1) = k_2 L_1 T_c + k_1 L_2 T_h$$

$$T_i = \frac{k_2 L_1 T_c + k_1 L_2 T_h}{k_1 L_2 + k_2 L_1}$$

$$Q = \frac{k_1 A}{L_1} (T_h - T_i) = \frac{k_1 A}{L_1 (k_1 L_2 + k_2 L_1)} (k_2 L_1 T_c + k_1 L_2 T_h - (k_1 L_2 + k_2 L_1) T_i)$$

recall  $\beta = KA \frac{T_h - T_c}{L}$

$K_{eff}$  (using both of  $L_1 + L_2$ )

$$K_{eff} = \frac{L_1 + L_2}{\left( L_1/k_1 + L_2/k_2 \right)}$$

mechanisms for heat conductivity

- diffusion [random molecular motions]
- convection [concentrated flows]
- radiation - emission of electromagnetic waves

Power emitted by a body is described by Stefan's law.

$$P = \sigma A e T^4$$

power radiated  $\uparrow$  emissivity

Stefan-Boltzmann constant

$$= 5.6696 \times 10^{-8} \frac{W}{m^2 K^4}$$

## § 21 Kinetic theory of gases

relate macroscopic

properties

to microscopic  
quantities

$P, V, T$

Here: low density gases.

- 1) # of molecules is large.  
separation is large.
- 2) molecules obey Newton's laws.  
but as a whole they move randomly  
'molecular randomness'