

1)(a) The probability distribution is $F(x) = Ce^{-\alpha x}$.
 First NORMALIZE. Note that the distance only takes
 on positive values.

$$\int_0^{\infty} F(x) dx = 1 = \int_0^{\infty} Ce^{-\alpha x} dx \quad y = \alpha x \quad dx = \frac{dy}{\alpha}$$

$$1 = C \int_0^{\infty} e^{-y} \frac{dy}{\alpha} = \frac{C}{\alpha} (1) = 1 \quad \text{so } \boxed{C = \alpha}$$

$$\lambda = \langle x \rangle = \int_0^{\infty} x Ce^{-\alpha x} dx = \alpha \int_0^{\infty} \left(\frac{y}{\alpha}\right) e^{-y} \left(\frac{dy}{\alpha}\right)$$

$$= \frac{1}{\alpha} \int_0^{\infty} ye^{-y} dy = \frac{1}{\alpha} \quad \boxed{\lambda = \frac{1}{\alpha}}$$

$$(b) P(x > 2\lambda) = \int_{2\lambda}^{\infty} F(x) dx = C \int_{2\lambda}^{\infty} e^{-\alpha x} dx$$

$$= \alpha \left[-\frac{1}{\alpha} e^{-\alpha x} \right]_{2\lambda}^{\infty} = \alpha \cdot \frac{1}{\alpha} e^{-2\lambda\alpha} = \boxed{e^{-2}}$$

2) The assumptions are i) circular orbits with $F = m \frac{v^2}{r}$
 ii) $L = mrv = nh$

$$F_0 = \frac{mv^2}{r} \quad * \quad v = \frac{nh}{mr}$$

$$F_0 = \frac{m}{r} \left(\frac{nh}{mr}\right)^2 = \frac{n^2 h^2}{mr^3} \quad r^3 = \frac{n^2 h^2}{mF_0}$$

$$\boxed{r_n = \left[\frac{n^2 h^2}{mF_0} \right]^{\frac{1}{3}}}$$

For the energy

$$E = \frac{1}{2} mv^2 + V = \frac{1}{2} mv^2 + F_0 r$$

but from above

$$mv^2 = F_0 r \quad \Rightarrow \quad E = \frac{3}{2} F_0 r$$

$$E_n = \frac{3}{2} F_0 \left[\frac{n^2 \hbar^2}{m F_0} \right]^{\frac{1}{3}} = \frac{3}{2} \left[n^2 \hbar^2 F_0^2 / m \right]^{\frac{1}{3}}$$

(b) For the state to be bound we need $r_n < a$

$$\left[\frac{n^2 \hbar^2}{m F_0} \right] < a^3 \quad n^2 < m F_0 a^3 / \hbar^2$$

$$n < \left[m c^2 F_0 a^3 / (\hbar c)^2 \right]^{\frac{1}{2}}$$

$$= \left[(5.11 \times 10^5 \text{ eV}) \left(\frac{1 \text{ eV}}{\text{nm}} \right) (8 \text{ nm})^3 / (197.33 \text{ eV} \cdot \text{nm})^2 \right]^{\frac{1}{2}}$$

$$n < 10.25 \quad \Rightarrow \quad n=10 \text{ is the last one}$$

10 states

$$3) \quad \psi = C e^{-ax^2}$$

$$V(x) = \frac{1}{2} k x^2$$

$$\frac{d\psi}{dx} = -2ax C e^{-ax^2};$$

$$\frac{d^2\psi}{dx^2} = -2a C e^{-ax^2} - 2ax(-2ax) e^{-ax^2} \cdot C$$

$$= [-2a + 4a^2 x^2] C e^{-ax^2}$$

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi + V(x) \psi$$

$$= -\frac{\hbar^2}{2m} [-2a + 4a^2 x^2] C e^{-ax^2} + \frac{1}{2} k x^2 C e^{-ax^2}$$

$$= \left[\frac{\hbar^2 a}{m} - 2 \frac{\hbar^2 a^2}{m} x^2 + \frac{1}{2} k x^2 \right] C e^{-ax^2}$$

We need this to be [constant] $\psi \Rightarrow$ we need.

$$2 \frac{\hbar^2 a^2}{m} = \frac{1}{2} k$$

$$a^2 = \frac{1}{4} k m / \hbar^2$$

$$a = \frac{\sqrt{k m}}{2 \hbar}$$

and then

$$E = \frac{\hbar^2 a}{m} = \frac{\hbar^2}{m} \frac{\sqrt{k m}}{2 \hbar}$$

$$E = \frac{1}{2} \hbar \sqrt{\frac{k}{m}}$$

4) According to Planck's hypothesis the only allowed energies for a mode of frequency ν are

$$E = 0, h\nu, 2h\nu, 3h\nu \dots, nh\nu, \dots$$

The probability of having energy E_n is given by the Boltzmann formula

$$P(E_n) = C e^{-E_n/kT} = C e^{-nh\nu/kT}$$

Here

$$h\nu = \frac{hc}{\lambda} = (1240 \text{ eV}\cdot\text{nm}) / 600 \text{ nm} = 2.07 \text{ eV}$$

$$kT = (8.617 \times 10^{-5} \text{ eV/K})(1500 \text{ K}) = 0.129 \text{ eV}$$

so

$$\frac{h\nu}{kT} = 16 \quad e^{-h\nu/kT} = e^{-16} = 1.14 \times 10^{-7}$$

$h\nu \gg kT$ so nearly all modes have $E=0$

$$P(E=h\nu) / P(E=0) = e^{-h\nu/kT} = 1.14 \times 10^{-7}$$

$$P(E=0) \approx 1$$

$$\Rightarrow \boxed{P(E=h\nu) \approx 1.14 \times 10^{-7}}$$