

23) We have $\lambda = \frac{c}{\sqrt{\omega^2 - \omega_0^2}}$ $\lambda = \frac{2\pi}{k}$ $v = \frac{\omega}{2\pi}$

so $\frac{2\pi}{k} = c \left[\left(\frac{\omega}{2\pi}\right)^2 - \left(\frac{\omega_0}{2\pi}\right)^2 \right]^{-\frac{1}{2}}$ $\omega_0 = 2\pi \nu_0$

so $k = 2\pi c \left[\left(\frac{\omega}{2\pi}\right)^2 - \left(\frac{\omega_0}{2\pi}\right)^2 \right]^{\frac{1}{2}} = c \left[\omega^2 - \omega_0^2 \right]^{\frac{1}{2}}$

$$k^2 = c^2 (\omega^2 - \omega_0^2) \Rightarrow \omega^2 = \omega_0^2 + \frac{k^2}{c^2}$$

$$\omega = \left[\omega_0^2 + \frac{k^2}{c^2} \right]^{\frac{1}{2}}$$

$$v_g = \frac{d\omega}{dk} = \frac{1}{2} \left[\omega_0^2 + \frac{k^2}{c^2} \right]^{-\frac{1}{2}} \left(\frac{2k}{c^2} \right) = \boxed{\frac{1}{c^2} \frac{k}{\omega}}$$

24) $\omega^2 = gk \frac{(1 - e^{-2kh})}{(1 + e^{-2kh})}$

a) When $h \gg \lambda$ then $h \gg \frac{2\pi}{k} \Rightarrow hk \gg 2\pi$

so for $h \gg \lambda$

$$e^{-2kh} \rightarrow e^{-\infty} \rightarrow 0 \Rightarrow \omega^2 = gk$$

$$\omega = \sqrt{gk}$$

$$\boxed{v_{\text{phase}} = \frac{\omega}{k} = \sqrt{\frac{g}{k}}}$$

$$v_{\text{group}} = \frac{d\omega}{dk} = \frac{d}{dk} \sqrt{g} k^{\frac{1}{2}} = \frac{1}{2} \sqrt{g} k^{-\frac{1}{2}} \Rightarrow \boxed{v_{\text{group}} = \frac{1}{2} \sqrt{\frac{g}{k}}}$$

b) For $h \ll \lambda$ $hk \rightarrow 0$ $e^x = 1 + x + \frac{x^2}{2} + \dots$

$$e^{-2hk} = 1 - 2hk + 2(hk)^2 - \dots$$

$$\omega^2 = gk \left[\frac{1 - (1 - 2hk + \dots)}{1 + (1 - 2hk + \dots)} \right] = gk \frac{2hk}{2} = ghk^2$$

$$\omega = \sqrt{gh} k \quad \boxed{v_p = \frac{\omega}{k} = \sqrt{gh} = \frac{d\omega}{dk} = v_g} \Rightarrow \text{non-dispersive}$$

25) So the wave packet is going to spread. I will assume that the "size" of the packet means σ_x at $t=0$. We can either use equations from the text or from class notes. I started with

$$A(k) = C e^{-a(k-k_0)^2}$$

and found a probability distribution

$$|\Psi(x,t)|^2 = \frac{|C|^2}{2} \left[\frac{1}{a^2 + \delta^2 t^2} \right]^{\frac{1}{2}} e^{-(x-\beta t)^2 / 2a(1 + \frac{\delta^2 t^2}{a^2})}$$

The "standard" gaussian formula is $e^{-u^2/2\sigma^2}$ so we have

$$\sigma = \sqrt{a} \left(1 + \frac{\delta^2 t^2}{a^2} \right)^{\frac{1}{2}}$$

At $t=0$ $\sigma = \sqrt{a} = 10^{-3} \text{ m}$ so $a = 10^{-6} \text{ m}^2$

The constant δ is defined as

$$\delta = \frac{1}{2} \frac{d^2 \omega}{dk^2} \Big|_{k_0}$$

For non-relativistic particles $E = p^2/2m \Rightarrow \omega = \frac{\hbar k^2}{2m}$

and

$$\delta = \frac{1}{2} \frac{d^2 \omega}{dk^2} = \frac{1}{2} \frac{d^2}{dk^2} \frac{\hbar k^2}{2m} = \frac{1}{2} \frac{d}{dk} \frac{\hbar k}{m} = \frac{\hbar}{2m}$$

\Rightarrow

$$\sigma = \sqrt{a} \left(1 + \left(\frac{\hbar t}{2ma} \right)^2 \right)^{\frac{1}{2}}$$

Lets find the transit time. $E = \frac{p^2}{2m} = 13.6 \text{ eV} \Rightarrow p = \sqrt{2mE}$

$$v = \frac{p}{m} = \sqrt{\frac{2E}{m}}$$

$$t = \frac{d}{v} = d \sqrt{\frac{m}{2E}}$$

$$\sigma = \sqrt{a} \left[1 + \frac{\hbar^2 d^2 \left(\frac{m}{2E} \right)}{4 m^2 a^2} \right]^{\frac{1}{2}} = \sqrt{a} \left[1 + \frac{\hbar^2 d^2}{8 m E a^2} \right]^{\frac{1}{2}}$$

$$\sigma = (10^{-3} \text{ m}) \left[1 + \frac{(197 \text{ eV} \cdot \text{nm})^2 \left(\frac{10^{-9} \text{ m}}{\text{nm}} \right)^2}{8 (5.11 \times 10^5 \text{ eV}) (13.6 \text{ eV}) \left(\frac{10^7 \text{ m}}{10^9 \text{ m}^2} \right)^2} \right]^{\frac{1}{2}}$$

$$= (10^{-3} \text{ m}) \left[1 + 7.0 \times 10^4 \right]^{\frac{1}{2}} = \boxed{0.265 \text{ m}}$$

For $KE = 100 \text{ MeV}$ the particle is relativistic so we need to re-evaluate δ .

$$E = [(pc)^2 + (mc^2)^2]^{\frac{1}{2}}$$

so

$$\omega = \frac{1}{\hbar} [(\hbar kc)^2 + (mc^2)^2]^{\frac{1}{2}}$$

$$\begin{aligned} \frac{d\omega}{dk} &= \frac{1}{\hbar} \left(\frac{1}{2}\right) [(\hbar kc)^2 + (mc^2)^2]^{-\frac{1}{2}} (2\hbar^2 kc^2) \\ &= \hbar kc^2 [(\hbar kc)^2 + (mc^2)^2]^{-\frac{1}{2}} \end{aligned}$$

$$\begin{aligned} \frac{d^2\omega}{dk^2} &= \hbar c^2 [(\hbar kc)^2 + (mc^2)^2]^{-\frac{1}{2}} - \frac{1}{2} \hbar kc^2 [\dots]^{-\frac{3}{2}} (2\hbar^2 kc^2) \\ &= \frac{\hbar c^2}{E} \left[1 - \frac{p^2 c^2}{E^2} \right] = \frac{\hbar c^2}{E} \left[\frac{E^2 - (pc)^2}{E^2} \right] \\ &= \frac{\hbar c^2}{E} \left[\frac{(mc^2)^2}{E^2} \right] = \hbar c^2 \frac{(mc^2)^2}{E^3} \end{aligned}$$

Here $KE = 100 \text{ MeV}$ and $mc^2 = 0.511 \text{ MeV} \Rightarrow E = 100.5 \text{ MeV}$.

We can get the velocity from δ . $E = \gamma mc^2 \Rightarrow$

$$\gamma = \frac{100.511}{0.511} = \frac{1}{\sqrt{1 - v^2/c^2}} \quad v = 0.999987c \approx c$$

$$\begin{aligned} \frac{\delta t}{a} &= \frac{1}{2} \frac{d^2\omega}{dk^2} \cdot \left(\frac{d}{a}\right) \frac{1}{v} = \\ &= \frac{1}{2} \hbar c^2 \frac{(mc^2)^2}{E^3} \left(\frac{d}{a}\right) \frac{1}{v} = \frac{1}{2} \hbar c \frac{(mc^2)^2}{E^3} \frac{d}{a} \frac{c}{v} \\ &= \frac{1}{2} (197.3 \text{ eV}\cdot\text{nm}) \frac{(5.11 \times 10^5 \text{ eV})^2}{(100.511 \times 10^6 \text{ eV})^3} \left(\frac{10^7 \text{ m}}{10^{-6} \text{ m}^2}\right) \left(\frac{1}{0.999987}\right) (10^{-9} \text{ m}) \\ &= 2.5 \times 10^{-7} \end{aligned}$$

So

$$\sigma = \sqrt{a} \left(1 + \left(\frac{\delta t}{a}\right)^2 \right)^{\frac{1}{2}} = \sqrt{a} \left(1 + (2.5 \times 10^{-7})^2 \right)^{\frac{1}{2}} = \boxed{10^{-3} \text{ m}}$$

26) (a) $\Psi = A e^{-\lambda|x|} e^{-i\omega t} \Rightarrow |\Psi|^2 = |A|^2 e^{-2\lambda|x|} = P(x)$
 Notice that $P(x)$ is an even function so the normalization condition is

$$\int_{-\infty}^{\infty} P(x) dx = 2 \int_0^{\infty} |A|^2 e^{-2\lambda x} dx = 2|A|^2 \left. \frac{1}{-2\lambda} e^{-2\lambda x} \right|_0^{\infty}$$

$$= 2|A|^2 \left(\frac{1}{2\lambda} \right) = \frac{|A|^2}{\lambda} = 1 \Rightarrow \boxed{A = \sqrt{\lambda}}$$

(b) $\langle x \rangle = \int_{-\infty}^{\infty} x P(x) dx = 0$ since the integrand is odd.

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 P(x) dx = 2 \int_0^{\infty} x^2 P(x) dx$$

$$= 2\lambda \int_0^{\infty} x^2 e^{-2\lambda x} dx$$

$$y = 2\lambda x$$

$$dx = \frac{1}{2\lambda} dy$$

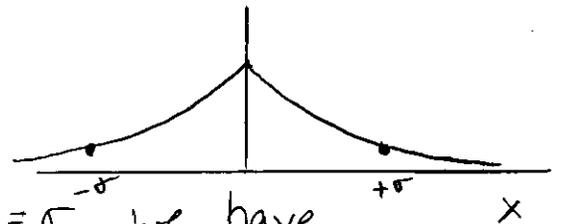
$$= 2\lambda \left(\frac{1}{2\lambda} \right)^3 \int_0^{\infty} y^2 e^{-y} dy$$

$$= \left(\frac{1}{2\lambda} \right)^2 \cdot 2$$

$$\boxed{\langle x^2 \rangle = \frac{1}{2\lambda^2}}$$

(c) $\sigma = \frac{1}{\sqrt{2}\lambda}$. The peak value of $|\Psi|^2$ is $|A|^2$

and at the points $x = \pm \sigma$, $|x| = \sigma$ we have



$$P(x) = |A|^2 e^{-2\lambda(\frac{1}{\sqrt{2}\lambda})} = |A|^2 e^{-\sqrt{2}} = 0.24 |A|^2$$

$$P_{\text{outside}} = 2 \int_{\sigma}^{\infty} P(x) dx = 2\lambda \int_{\sigma}^{\infty} e^{-2\lambda x} dx$$

$$= 2\lambda \left. \frac{1}{-2\lambda} e^{-2\lambda x} \right|_{\sigma}^{\infty} = e^{-2\lambda\sigma} = \boxed{0.24}$$

27)(a) We have

$$\psi = Ce^{\alpha x} + De^{-\alpha x}$$

$$\psi^* = C^*e^{\alpha x} + D^*e^{-\alpha x}$$

$$j = \frac{\hbar}{2im} \left[\psi^* \frac{d\psi}{dx} - \psi \frac{d\psi^*}{dx} \right]$$

$$= \frac{\hbar}{2im} \left[(C^*e^{\alpha x} + D^*e^{-\alpha x})(\alpha Ce^{\alpha x} - \alpha De^{-\alpha x}) \right. \\ \left. - (Ce^{\alpha x} + De^{-\alpha x})(\alpha C^*e^{\alpha x} - \alpha D^*e^{-\alpha x}) \right]$$

$$= \frac{\hbar\alpha}{2im} \left[\cancel{C^*C}e^{2\alpha x} - \cancel{D^*D}e^{-2\alpha x} + D^*C - C^*D \right. \\ \left. - \cancel{C}e^{2\alpha x} + \cancel{D}e^{-2\alpha x} + CD^* - DC^* \right]$$

$$j = \frac{\hbar\alpha}{im} (D^*C - C^*D)$$

(b) Matching Ae^{ikx} and $Ce^{\alpha x} + De^{-\alpha x}$ @ $x=a$.

$$\textcircled{1} Ae^{ika} = Ce^{\alpha a} + De^{-\alpha a}$$

$$\textcircled{2} ikAe^{ika} = \alpha Ce^{\alpha a} - \alpha De^{-\alpha a}$$

Multiply $\textcircled{1}$ by α , then add and subtract.

$$\text{ADD: } (ik + \alpha)Ae^{ika} = 2\alpha Ce^{\alpha a} \quad C = \frac{(ik + \alpha)}{2\alpha} Ae^{ika} e^{-\alpha a}$$

$$\text{SUBT: } (\alpha - ik)Ae^{ika} = 2\alpha De^{-\alpha a} \quad D = \frac{\alpha - ik}{2\alpha} Ae^{ika} e^{\alpha a}$$

Then starting from j of part (a)

$$j = \frac{\hbar\alpha}{im} \left[\left(\frac{\alpha + ik}{2\alpha} A^* e^{-ika} e^{\alpha a} \right) \left(\frac{\alpha + ik}{2\alpha} Ae^{ika} e^{-\alpha a} \right) \right. \\ \left. - \left(\frac{\alpha - ik}{2\alpha} A^* e^{-ika} e^{-\alpha a} \right) \left(\frac{\alpha - ik}{2\alpha} Ae^{ika} e^{\alpha a} \right) \right]$$

$$\begin{aligned}
 j &= \frac{\hbar d}{im} A^* A \left[\left(\frac{d+ik}{2d} \right)^2 - \left(\frac{d-ik}{2d} \right)^2 \right] \\
 &= \frac{\hbar d}{im} A^* A \left[\frac{d^2 + 2ikd - k^2 - (d^2 - 2ikd - k^2)}{4d^2} \right] \\
 &= \frac{\hbar d}{im} |A|^2 \left[\frac{4ikd}{4d^2} \right] = \frac{\hbar d}{im} |A|^2 \left(\frac{ik}{d} \right) = \frac{\hbar k}{m} |A|^2
 \end{aligned}$$

which is the correct j for $\psi = Ae^{ikx}$

28) (a) For $x < 0$ and $x > a$ $V = 0$ so we have

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi = E\psi \quad \frac{d^2}{dx^2} \psi = -k^2 \psi \quad k = \left[\frac{2mE}{\hbar^2} \right]^{\frac{1}{2}}$$

For $x < 0$ write

$$\psi_I = Ae^{ikx} + Be^{-ikx}$$

For $x > a$ "

$$\psi_{III} = Fe^{ikx} + Ge^{-ikx} = Fe^{ikx}$$

As usual we set $G = 0$ since we expect no left-traveling particles for $x > a$.

For $0 < x < a$ $V = -V_0$ so

$$\begin{aligned}
 -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi - V_0 \psi &= E\psi & \frac{d^2}{dx^2} \psi &= -\frac{2m}{\hbar^2} (E + V_0) \psi \\
 & & &= -g^2 \psi
 \end{aligned}$$

where $g = \left[\frac{2m(E + V_0)}{\hbar^2} \right]^{\frac{1}{2}}$

$$\psi_{II} = Ce^{igx} + De^{-igx}$$

MATCHING:

$$x = 0 \quad \psi \quad A + B = C + D \quad \textcircled{1}$$

$$\psi' \quad ik(A - B) = ig(C - D) \quad \textcircled{2}$$

$$x = a$$

$$\psi \quad C e^{iga} + D e^{-iga} = F e^{ika} \quad (3)$$

$$\psi' \quad ig(C e^{iga} - D e^{-iga}) = ik F e^{ika} \quad (4)$$

Algebra: Multiply (3) by g , then add and subtract (3) and (4)

$$(5) \quad (g+k)A + (g-k)B = 2gC \quad (6) \quad (g-k)A + (g+k)B = 2gD$$

Do the same with (3) and (4)

$$(7) \quad 2gC e^{iga} = (g+k) F e^{ika} \quad (8) \quad 2gD e^{-iga} = (g-k) F e^{ika}$$

Eliminate C between (5) and (7) and D between (6) and (8)

$$[(g+k)A + (g-k)B] e^{iga} = (g+k) F e^{ika} \quad (9)$$

$$[(g-k)A + (g+k)B] e^{-iga} = (g-k) F e^{ika} \quad (10)$$

Finally eliminate B . Multiply (9) by $(g+k)e^{-iga}$ and (10) by $(g-k)e^{iga}$ and subtract

$$[(g+k)^2 - (g-k)^2] A = [(g+k)^2 e^{-iga} - (g-k)^2 e^{iga}] F e^{ika}$$

$$\frac{F}{A} = \frac{(g+k)^2 - (g-k)^2}{[(g+k)^2 e^{-iga} - (g-k)^2 e^{iga}]} e^{-ika}$$

$$= \frac{(g+k)^2 - (g-k)^2}{[(g+k)^2 - (g-k)^2 e^{2iga}]} e^{-ika} e^{iga}$$

$$(c) \quad T = j_{\text{trans}} / j_{\text{inc}} = |F/A|^2$$

$$\frac{[(g+k)^2 - (g-k)^2]^2}{[(g+k)^2 - (g-k)^2 e^{2iga}] \cdot [(g+k)^2 - (g-k)^2 e^{-2iga}]}$$

$$T = \frac{(g+k)^4 + (g-k)^4 - 2(g+k)^2(g-k)^2}{(g+k)^4 + (g-k)^4 - 2(g+k)^2(g-k)^2 \cos 2ga}$$

We get $T=1$ whenever $\cos 2ga = 1 \Rightarrow$

$$2ga = 0, 2\pi, 4\pi \dots$$

$$a_1 = \frac{2\pi}{2g} = \frac{\pi}{g}$$

$$g = \left[\frac{2m}{\hbar^2} (V_0 + E) \right]^{\frac{1}{2}} = \left[2mc^2 (V_0 + E) \right]^{\frac{1}{2}} / \hbar c$$

$$= \left[2 (5.11 \times 10^5 \text{ eV}) (8 \text{ eV}) \right]^{\frac{1}{2}} / 197.3 \text{ eV} \cdot \text{nm} = 14.5 / \text{nm}$$

$$\boxed{a = 0.217 \text{ nm}}$$