

## HOMEWORK 4 SOLUTIONS

15) So we need to calculate  $\langle \phi_{1s}(1) \phi_{2s}(2) | \frac{e^2}{4\pi\epsilon_0 |\vec{r}_1 - \vec{r}_2|} | \phi_{2s}(1) \phi_{1s}(2) \rangle$ .

where

$$\phi_{1s} = \frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} e^{-zr/a_0}$$

$$\phi_{2s} = \frac{1}{\sqrt{\pi}} \left(\frac{z}{2a_0}\right)^{3/2} \left(1 - \frac{zr}{2a_0}\right) e^{-zr/2a_0}$$

So when we put everything together we have

$$\Delta E = \frac{e^2}{4\pi\epsilon_0} \iint \rho_1(\vec{r}_1) \rho_2(\vec{r}_2) \frac{1}{|\vec{r}_1 - \vec{r}_2|} d^3r_1 d^3r_2$$

where  $\rho_1 = \rho_2 = \phi_{1s}(r) \phi_{2s}(r) = \frac{1}{\pi} \left(\frac{z}{a_0}\right)^3 \left(\frac{1}{2}\right)^{3/2} \left(1 - \frac{zr}{2a_0}\right) e^{-3zr/2a_0}$

We divide space into regions with  $r_1 > r_2$  and  $r_1 < r_2$ .

Then do the angular integrals ( $4\pi$ ) for  $d\Omega_1$  and ( $4\pi$ ) for  $d\Omega_2$ . Then do the Gauss's Law trick to get

$$\Delta E = \frac{e^2}{4\pi\epsilon_0} (4\pi)^2 \left\{ \int_0^\infty \rho_2(r_2) \frac{1}{r_2} \left[ \int_0^{r_2} \rho_1(r_1) r_1^2 dr_1 \right] r_2^2 dr_2 + \text{a second identical term} \right\}$$

$$= \left(\frac{e^2}{4\pi\epsilon_0}\right) (4\pi)^2 \left(\frac{1}{\pi}\right)^2 \left(\frac{z}{a_0}\right)^6 \left(\frac{1}{8}\right) (2) \int_0^\infty \left(1 - \frac{zr_2}{2a_0}\right) e^{-3zr_2/2a_0} \times \frac{1}{r_2} \left[ \int_0^{r_2} \left(1 - \frac{zr_1}{2a_0}\right) e^{-3zr_1/2a_0} r_1^2 dr_1 \right] r_2^2 dr_2$$

Lets do the inner integral. Let

$$x_1 = \frac{3z}{2a_0} r_1$$

$$x_2 = \frac{3z}{2a_0} r_2$$

$$[\ ] = \int_0^{x_2} \left(1 - \frac{1}{3} x_1\right) e^{-x_1} \left(\frac{2a_0}{3z}\right)^3 x_1^2 dx_1$$

$$= \left(\frac{2a_0}{3z}\right)^3 \left\{ -\left(x_1^2 + 2x_1 + 2\right) e^{-x_1} + \frac{1}{3} \left(x_1^3 + 3x_1^2 + 6x_1 + 6\right) e^{-x_1} \right\} \Big|_0^{x_2}$$

$$= \left(\frac{2a_0}{3z}\right)^3 \left\{ \frac{1}{3} x_2^3 e^{-x_2} \right\}$$

So

$$\begin{aligned} \Delta E &= 4 \left( \frac{e^2}{4\pi\epsilon_0} \right) \left( \frac{Z}{a_0} \right)^6 \int_0^\infty \left( 1 - \frac{Zr_2}{2a_0} \right) e^{-3Zr_2/2a_0} \frac{1}{r_2} \\ &\quad \times \left( \frac{2a_0}{3Z} \right)^3 \left( \frac{1}{3} \right) \left( \frac{3Z}{2a_0} r_2 \right)^3 e^{-3Zr_2/2a_0} r_2^2 dr_2 \\ &= \frac{4}{3} \left( \frac{e^2}{4\pi\epsilon_0} \right) \left( \frac{Z}{a_0} \right)^6 \int_0^\infty \left( 1 - \frac{Zr_2}{2a_0} \right) r_2^4 e^{-3Zr_2/2a_0} dr_2 \end{aligned}$$

Let

$$x = \frac{3Z}{2a_0} r_2 \Rightarrow$$

$$\begin{aligned} \Delta E &= \left( \frac{4}{3} \right) \left( \frac{e^2}{4\pi\epsilon_0} \right) \left( \frac{Z}{a_0} \right)^6 \left( \frac{2a_0}{3Z} \right)^5 \int_0^\infty \left( 1 - \frac{x}{6} \right) x^4 e^{-x} dx \\ &= \frac{4}{3} \left( \frac{e^2}{4\pi\epsilon_0} \right) \left( \frac{Z}{a_0} \right) \left( \frac{1}{3} \right)^5 [4! - \frac{1}{6} \cdot 5!] \\ &= \frac{4}{3^6} \left( \frac{e^2}{4\pi\epsilon_0} \right) \frac{Z}{a_0} [24 - 20] = \boxed{\frac{2^4}{3^6} \frac{e^2}{4\pi\epsilon_0} \frac{Z}{a_0}} \end{aligned}$$

16) For  $(1s)'(2s)'$  the starting energy (neglecting  $V_{ee}$ ) is

$$\begin{aligned} E^{(0)} &= -\frac{mc^2}{2} \alpha^2 \left[ \frac{Z^2}{(1)^2} + \frac{(Z)^2}{(2)^2} \right] = -[54.4 \text{ eV} + 13.6 \text{ eV}] \\ &= -68.0 \text{ eV} \end{aligned}$$

Then the corrections are

$$\begin{aligned} \Delta E^{(1)} &= \Delta E_{\text{DIR}} \pm \Delta E_{\text{EX}} \\ &= \left[ \left( \frac{17}{81} \right) \pm \frac{2^4}{3^6} \right] \frac{e^2}{4\pi\epsilon_0} \frac{Z}{a_0} = [ ] Z mc^2 \alpha^2 \\ &= 11.42 \text{ eV} \pm 1.19 \text{ eV} \end{aligned}$$

$\Rightarrow$  we predict

$$E(3s) = -68.0 + 11.42 - 1.19 = -57.77 \text{ eV}$$

$$E(1s) = \quad \quad \quad \quad + 1.19 = -55.39 \text{ eV.}$$

$$17) \quad H = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \quad \text{where } V(x) = \begin{cases} \infty & x < 0 \\ Cx & x > 0 \end{cases}$$

$$\psi_t = Nxe^{-\alpha x} \quad \text{for } x > 0$$

Normalize

$$\langle \psi_t | \psi_t \rangle = N^2 \int_0^{\infty} x^2 e^{-2\alpha x} dx$$

$$\text{Let } y = 2\alpha x \quad dx = y/2\alpha$$

$$\langle \psi_t | \psi_t \rangle = \frac{N^2}{(2\alpha)^3} \int_0^{\infty} y^2 e^{-y} dy = \frac{N^2}{(2\alpha)^3} \cdot 2 = 1$$

$$\boxed{N^2 = 4\alpha^3}$$

$$E_t = \langle \psi_t | H | \psi_t \rangle$$

$$\frac{d^2}{dx^2} \psi = N \frac{d}{dx} (e^{-\alpha x} - \alpha x e^{-\alpha x}) = N [-\alpha e^{-\alpha x} - \alpha e^{-\alpha x} + \alpha^2 x e^{-\alpha x}]$$

$$= \alpha N [\alpha x - 2] e^{-\alpha x}$$

So

$$E_t = -\frac{\hbar^2}{2m} \alpha N^2 \int_0^{\infty} x e^{-\alpha x} (\alpha x - 2) e^{-\alpha x} dx$$

$$+ CN^2 \int_0^{\infty} x e^{-\alpha x} \cdot x \cdot x e^{-\alpha x} dx$$

$$= N^2 \left[ -\frac{\hbar^2 \alpha}{2m} \int_0^{\infty} (\alpha x^2 - 2x) e^{-2\alpha x} dx \right.$$

$$\left. + C \int_0^{\infty} x^3 e^{-2\alpha x} dx \right]$$

$$\text{Let } y = 2\alpha x \Rightarrow$$

$$E_t = N^2 \left[ -\frac{\hbar^2 \alpha}{2m} \int_0^{\infty} \left[ \frac{y^2}{4\alpha} - 2 \frac{y}{2\alpha} \right] e^{-y} \frac{dy}{2\alpha} \right.$$

$$\left. + C \int_0^{\infty} \left( \frac{1}{2\alpha} \right)^4 y^3 e^{-y} dy \right]$$

$$E_t = N^2 \left[ -\frac{\hbar^2 \alpha}{2m} \left( \frac{1}{2\alpha} \right) \left( \frac{1}{4\alpha} \cdot 2 - \frac{1}{\alpha} \right) + C \left( \frac{1}{2\alpha} \right)^4 \cdot 6 \right]$$

$$\Rightarrow E_t = (4\alpha^3) \left[ +\frac{\hbar^2}{2m} \left( \frac{1}{4\alpha^2} \right) + \frac{3}{8} \frac{C}{\alpha^4} \right]$$

$$= \frac{\hbar^2 \alpha^2}{2m} + \frac{3}{2} \frac{C}{\alpha}$$

Now choose  $\alpha$  to minimize  $E_t$ .

$$\frac{dE_t}{d\alpha} = 2 \frac{\hbar^2 \alpha}{2m} - \frac{3}{2} \frac{C}{\alpha^2} = 0 \quad \frac{\hbar^2 \alpha^3}{m} = \frac{3}{2} C$$

$$\alpha^3 = \frac{3}{2} \frac{Cm}{\hbar^2} \quad \alpha = \left( \frac{3}{2} \frac{Cm}{\hbar^2} \right)^{\frac{1}{3}}$$

which gives

$$E_t = \frac{\hbar^2}{2m} \left( \frac{3}{2} \frac{Cm}{\hbar^2} \right)^{\frac{2}{3}} + \frac{3}{2} C \left( \frac{2}{3} \frac{\hbar^2}{Cm} \right)^{\frac{1}{3}}$$

Lets combine the two terms

$$E_t = \frac{\hbar^2}{2m} \frac{3^{2/3}}{2^{2/3}} \frac{C^{2/3} m^{2/3}}{\hbar^{4/3}} + \frac{3}{2} C \left( \frac{2^{1/3}}{3^{1/3}} \right) \frac{\hbar^{2/3}}{C^{1/3} m^{1/3}}$$

$$= \frac{1}{2} \left( \frac{3}{2} \right)^{\frac{2}{3}} \frac{\hbar^{2/3} C^{2/3}}{m^{1/3}} + \left( \frac{3}{2} \right)^{\frac{2}{3}} \frac{\hbar^{2/3} C^{2/3}}{m^{1/3}}$$

$$= \left( \frac{3}{2} \right) \left( \frac{3}{2} \right)^{\frac{2}{3}} \frac{\hbar^{2/3} C^{2/3}}{m^{1/3}} = \boxed{\frac{3}{2} \left[ \frac{9}{4} \frac{\hbar^2 C^2}{m} \right]^{\frac{1}{3}}}$$

18) The helium ground state energy is  $-78.975 \text{ eV}$ . According to the graph the wavelength to excite the  $(2s)'(2p)'$  state is about  $203.7 \text{ \AA} = 20.37 \text{ nm}$  corresponding to

$$E = \frac{hc}{\lambda} = \frac{1240 \text{ eV}\cdot\text{nm}}{20.37 \text{ nm}} = 60.87 \text{ eV}$$

and so the state is at  $E = -79.975 + 60.87 = -18.10 \text{ eV}$ .  
 If the atom autoionizes to the ground state of  $\text{He}^+$  the energy of the "atom" is  $-54.4 \text{ eV}$  and so the emitted electron has (by energy conservation)

$$E = 54.4 \text{ eV} - 18.10 \text{ eV} = 36.3 \text{ eV}$$

$\Rightarrow$

$$\frac{1}{2} m v^2 = 36.3 \text{ eV} = E \quad v = [2E/m]^{\frac{1}{2}} = [2E/mc^2]^{\frac{1}{2}} \cdot c$$

$$v = \left[ \frac{2(36.3 \text{ eV})}{5.11 \times 10^5 \text{ eV}} \right]^{\frac{1}{2}} \cdot c = 0.012c = \boxed{3.57 \times 10^6 \text{ m/s}}$$

If  $\text{He}^+$  is in its 1<sup>st</sup> excited state the energy would be  $-54.4 \text{ eV}/4 = -13.6 \text{ eV}$ . This would require the emitted electron to have negative KE so this decay does not occur.

19)  $H = \frac{p^2}{2m} + \lambda x^4$

We try

$$\psi_t = N e^{-\alpha x^2}$$

Normalize

$$\langle \psi_t | \psi_t \rangle = N^2 \int_{-\infty}^{\infty} e^{-2\alpha x^2} dx$$

$$= 2N^2 \int_0^{\infty} e^{-2\alpha x^2} dx$$

Using integral 666

$$\int_0^{\infty} x^{2n} e^{-ax^2} dx = \frac{1 \cdot 3 \cdot 5 \cdots (2n-1)}{2^{n+1} a^n} \sqrt{\frac{\pi}{a}}$$

$$\langle \psi_t | \psi_t \rangle = \frac{1}{2} \sqrt{\frac{\pi}{2\alpha}} \cdot 2N^2 = N^2 \sqrt{\frac{\pi}{2\alpha}} = 1 \Rightarrow N^2 = \sqrt{\frac{2\alpha}{\pi}}$$

$$\begin{aligned} \langle V \rangle &= \langle \psi_t | V | \psi_t \rangle = \lambda \int_{-\infty}^{\infty} \psi_t^* x^4 \psi_t dx \\ &= \lambda N^2 (2) \int_0^{\infty} x^4 e^{-2\alpha x^2} dx \\ &= \lambda N^2 (2) \frac{1.3}{8 (2\alpha)^2} \sqrt{\frac{\pi}{2\alpha}} \\ &= 2 \lambda \left(\frac{2\alpha}{\pi}\right)^{\frac{1}{2}} \frac{3}{32\alpha^2} \sqrt{\frac{\pi}{2\alpha}} = \frac{3}{16} \frac{\lambda}{\alpha^2} \end{aligned}$$

$$\begin{aligned} \frac{p^2}{2m} \psi_t &= -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} N e^{-\alpha x^2} = -\frac{\hbar^2}{2m} N \frac{d}{dx} (-2\alpha x e^{-\alpha x^2}) \\ &= -\frac{\hbar^2}{2m} N [-2\alpha e^{-\alpha x^2} + 4\alpha^2 x^2 e^{-\alpha x^2}] \end{aligned}$$

$$\begin{aligned} \Rightarrow \langle T \rangle &= -\frac{\hbar^2}{2m} N^2 \int_{-\infty}^{\infty} e^{-\alpha x^2} (-2\alpha + 4\alpha^2 x^2) e^{-\alpha x^2} dx \\ &= -\frac{\hbar^2}{2m} N^2 (2) \int_0^{\infty} -2\alpha e^{-2\alpha x^2} + 4\alpha^2 x^2 e^{-2\alpha x^2} dx \\ &= -\frac{\hbar^2}{2m} \sqrt{\frac{2\alpha}{\pi}} (2) \left[ -2\alpha \left(\frac{1}{2}\right) \sqrt{\frac{\pi}{2\alpha}} + 4\alpha^2 \frac{1}{4(2\alpha)} \sqrt{\frac{\pi}{2\alpha}} \right] \\ &= -\frac{\hbar^2}{m} \left[ -\alpha + \frac{\alpha}{2} \right] = +\frac{\hbar^2 \alpha}{2m} \end{aligned}$$

$$E_t = \langle T \rangle + \langle V \rangle = \frac{\hbar^2 \alpha}{2m} + \frac{3}{16} \frac{\lambda}{\alpha^2}$$

$$\frac{dE_t}{d\alpha} = \frac{\hbar^2}{2m} - \frac{3}{8} \frac{\lambda}{\alpha^3} = 0 \quad \frac{\hbar^2}{2m} = \frac{3}{8} \frac{\lambda}{\alpha^3}$$

$$\alpha^3 = \frac{3}{4} \frac{\lambda m}{\hbar^2} \quad \alpha = \left(\frac{3}{4}\right)^{\frac{1}{3}} \left(\frac{\lambda m}{\hbar^2}\right)^{\frac{1}{3}}$$

and so

$$E_t = \frac{\hbar^2}{2m} \left(\frac{3}{4}\right)^{\frac{1}{3}} \left(\frac{\lambda m}{\hbar^2}\right)^{\frac{1}{3}} + \frac{3}{16} \lambda \left(\frac{4}{3}\right)^{\frac{2}{3}} \left(\frac{\hbar^2}{\lambda m}\right)^{\frac{2}{3}}$$

$$E_t = \frac{1}{2} \left( \frac{3}{4} \right)^{\frac{1}{3}} \frac{\hbar^4 \lambda^{\frac{1}{3}}}{m^{\frac{2}{3}}} + \frac{3}{16} \left( \frac{4}{3} \right)^{\frac{2}{3}} \frac{\hbar^4 \lambda^{\frac{1}{3}}}{m^{\frac{2}{3}}}$$

$$= \left( \frac{1}{2} + \frac{1}{4} \right) \left( \frac{3}{4} \right)^{\frac{1}{3}} \left( \frac{\hbar^4 \lambda}{m^2} \right)^{\frac{1}{3}} = \left( \frac{3}{4} \right)^{\frac{4}{3}} \left( \frac{\hbar^4 \lambda}{m^2} \right)^{\frac{1}{3}}$$

$$= \left[ \frac{(3)^{\frac{4}{3}}}{4} \right] \lambda^{\frac{1}{3}} \left( \frac{\hbar^2}{2m} \right)^{\frac{2}{3}} = 1.0817 \lambda^{\frac{1}{3}} \left( \frac{\hbar^2}{2m} \right)^{\frac{2}{3}}$$

So our trial wave function gives an energy too high by around 2%